

Séminaire MLMS

ϕ -FEM :

A fictitious domain method for finite element methods
on domains defined by level-sets

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Outline

- The "standard" Finite Element Method
- Geometrical constraint on the mesh
- Previous Fictitious Domain Methods
- ϕ -FEM
- Conclusions and perspectives

Weak formulation

Strong formulation of the Poisson equation

$$\left\{ \begin{array}{l} \text{Find } u \in H^2(\Omega) \text{ s.t. :} \\ -\Delta u = f, \text{ in } \Omega \\ u = 0, \text{ on } \partial\Omega \end{array} \right.$$



$H^2(\Omega)$ \approx differentiable two times

Multiplying by a "test" function v & by integration by part, one has

Weak formulation

$$\left\{ \begin{array}{l} \text{Find } u \in H_0^1(\Omega) \text{ s.t. :} \\ \Leftrightarrow \int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} fv \quad \forall v \in H_0^1(\Omega). \end{array} \right.$$

$H_0^1(\Omega)$ \approx differentiable one times, equal to zero on the boundary

Standard FEM formulation

Weak formulation

$$\begin{cases} \text{Find } u \in H_0^1(\Omega) \text{ s.t. :} \\ \Leftrightarrow \int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} f v \quad \forall v \in H_0^1(\Omega). \end{cases}$$

Standard FEM formulation

$$\begin{cases} \text{Find } u_h \in V_h \text{ s.t. :} \\ \Leftrightarrow \int_{\Omega} \nabla u_h \cdot \nabla v_h = \int_{\Omega} f v_h \quad \forall v \in V_h. \end{cases}$$

where V_h is a subspace of $H_0^1(\Omega)$ of **finite dimensional**.

Standard FEM formulation

Finite element space

Let $V_h = \langle \psi_k \in H_0^1(\Omega) : k \in \{1, \dots, N\} \rangle \subset H_0^1(\Omega)$.

Equivalence with a matrix system

$$\left\{ \begin{array}{l} \text{Find } u_h = \sum U_{hk} \psi_k \in V_h \text{ s.t. :} \\ \int_{\Omega} \nabla u_h \cdot \nabla \psi_k = \int_{\Omega} f \psi_k \quad \forall k \end{array} \right. \Leftrightarrow \left\{ \begin{array}{l} \text{Find } U_h \in \mathbb{R}^N \text{ s.t. :} \\ A_h U_h = F_h \end{array} \right.$$

where

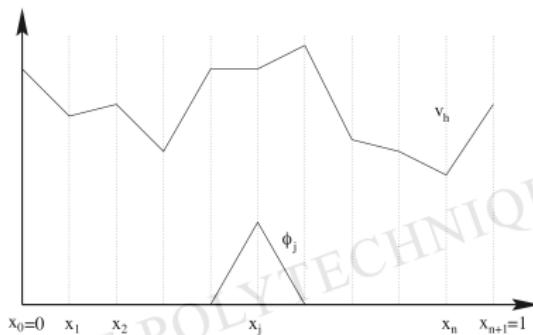
$$\left\{ \begin{array}{l} A_h = (\int_{\Omega} \nabla \psi_k \cdot \nabla \psi_j)_{kj} \\ F_h = (\int_{\Omega} f \psi_k)_k \\ U_h = (U_{hk})_k \end{array} \right.$$

Final solution

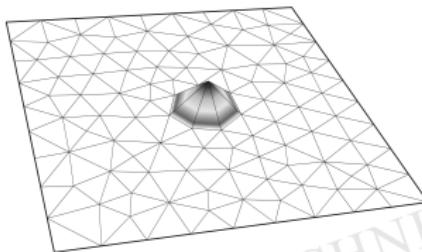
$$u_h = \sum U_{hk} \psi_k.$$

Lagrange continuous finite element

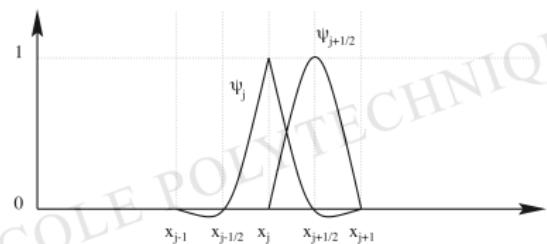
FE space : $V_h = \{\text{cont. piecewise pol. functions on a regular mesh}\}$



Piecewise linear Lagrange FE order 1 (\mathbb{P}_1)



Order 1 (\mathbb{P}_1), dimension 2



Order 2 (\mathbb{P}_2), dimension 1

Why use polynomials ?

because the computation of the coefficient $\int_{\Omega} \nabla \psi_k \cdot \nabla \psi_l$ of the FE matrix is explicit and **exact**

Important : if ψ_k is not polynomial, this computation is not exact anymore

Why a mesh ?

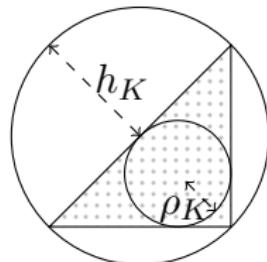
because we can **not good approximate a function** by a polynomial on the whole domain

Regular mesh : Geometrical condition

- **Simplex (triangle, tetrahedron)**

The standard FEM works under the **Ciarlet condition** [Ciarlet 78']

$$\frac{h_K}{\rho_K} < \gamma$$



- **In general (hexahedron,...)**

The standard FEM works if the Jacobian matrix to an reference element is positive.

Important remark : If the mesh contains degenerated cells :

- It is not guaranty that standard FEM converges
- The **conditioning number** of the FE matrix is bad

What is the conditioning number ?

Conditioning number of a matrix A :

$$\text{Cond}(\mathbf{A}) := \|\mathbf{A}\|_2 \|\mathbf{A}^{-1}\|_2$$

For symmetric definite positive matrices :

$$\text{Cond}(\mathbf{A}) := \frac{\text{bigest eigen value}}{\text{smallest eigen value}} > 1$$

Example

if $\mathbf{A} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \varepsilon \end{pmatrix}$ then $\text{Cond}(\mathbf{A}) = 1/\varepsilon$.

If ε goes to zero, \mathbf{A} goes to a non-invertible matrix

Conclusion

the conditioning number measures the **invertibility of the matrix**.

Poor conditioning of the system matrix

Proposition

- **Non-degenerated mesh :** Suppose that the mesh satisfies the geometrical conditions, then

$$\text{Cond}(\mathbf{A}) \leq C/h^2$$

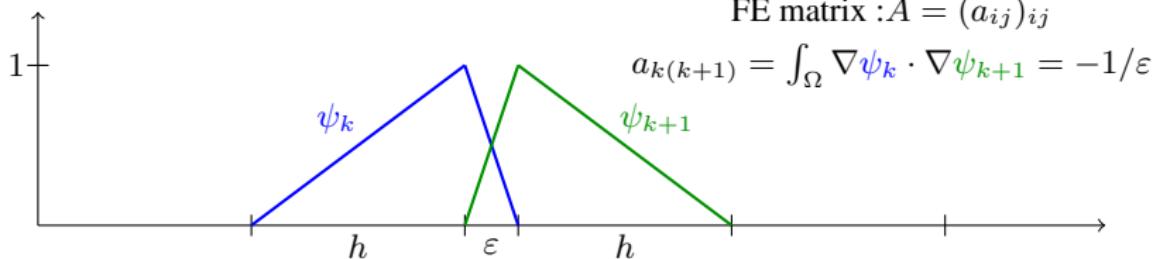
where h is the size of the cells.

- **Only one degenerated cell :**

$$\text{Cond}(\mathbf{A}) \geq C/h\varepsilon$$

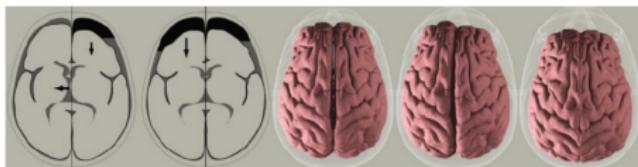
where ε is the size the degenerated cell.

Idea :



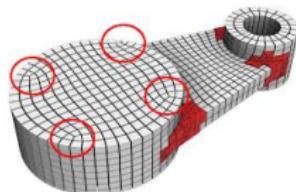
Complex geometries

What can we do on complex geometries ?



In particular how to use **hexahedral meshes**

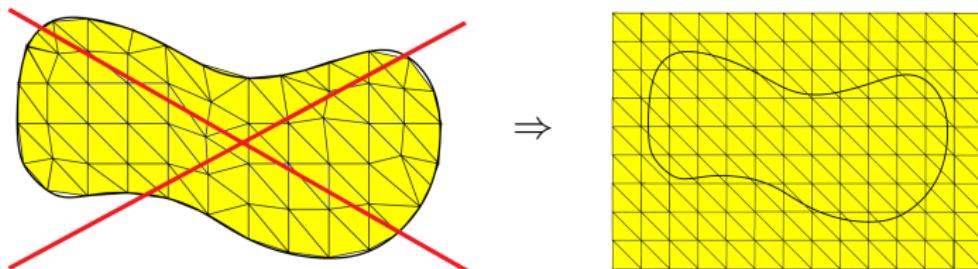
➤ Quasi-compressible elastic models : locking effect



What is the locking effect ?

- Theoretical continuous quasi-compressible models "converge" to theoretical continuous incompressible models
- False for numerical models with tetrahedrons !
 - not enough degree of freedom

Alternative : Fictitious domain methods

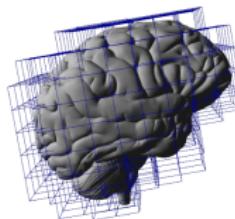


Advantages

- No need to mesh
- Regular cells

Difficulties

- Adapt the **weak formulation**
- **Conditioning** of the FE matrix



Outline

- The "standard" Finite Element Method
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- **Previous Fictitious domain Methods**
- ϕ -FEM
- Conclusions and perspectives

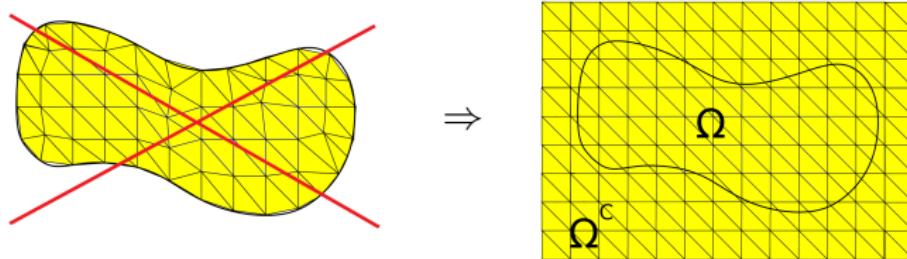
Previous works, first works

References

Saul'ev 63' (Dirichlet), Astrakhantsev 78' (Neumann),
Glowinski 92' (first proof)

Formal idea

$$\text{Total energy} = (\text{energy on } \Omega) + \varepsilon \times (\text{energy on } \Omega^c) \quad (\varepsilon \ll 1)$$



Advantage

Non-conform mesh (complex and time varying geometry)

Difficulty

Discontinuity, large FE matrix and bad cond. number

Previous works, XFEM : cut shape functions

References

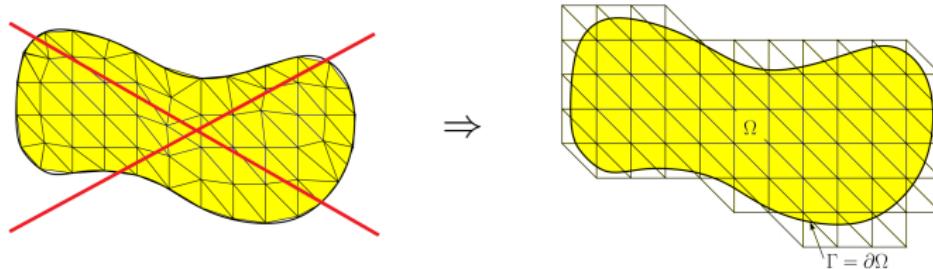
Moes-Béchet-Tourbier 2006, Haslinger-Renard 2009

Formal idea

Cut shape function : $\psi_k \rightarrow \psi_k \mathbb{1}_\Omega$

Boundary condition on Γ : Lagrange multiplier

Conditioning of the matrix : stabilization on the boundary



Advantage

Small FE matrix

Good conditioning number

Difficulty

Non-classical shape functions and **discontinuity** in the integrals

Previous works, Cut FEM : partial integration on the cells

References

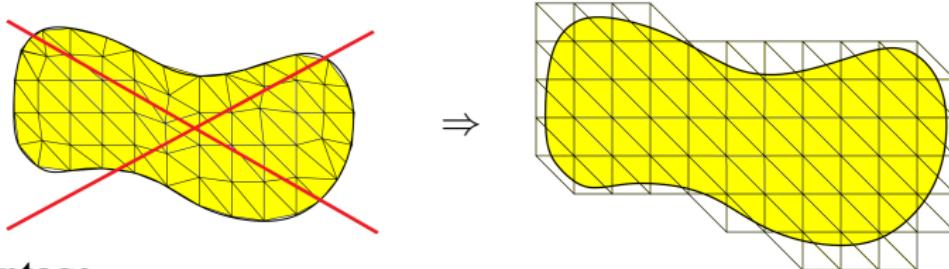
Burman et al 2010-2014

Formal idea

Partial integration on the cell near the boundary

Lagrange multiplier or penalization for the boundary conditions

Conditioning of the matrix : stabilization on the boundary



Advantage

Standard shape functions

Difficulty

- **Integral on the real boundary, cut integral.**
- \mathbb{P}_k formulation

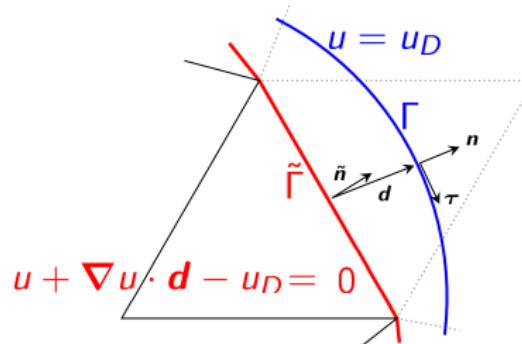
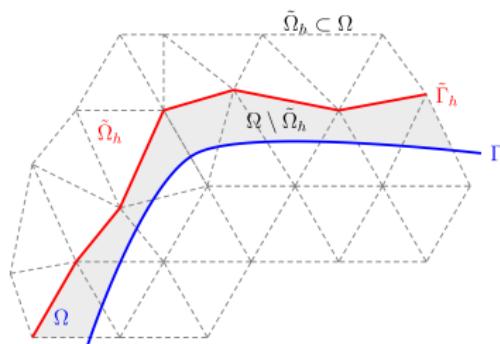
Previous work : Shifted Bound. Method : Taylor expansion

Formal idea (Dirichlet)

Taylor expansion of the bound. cond. Main-Scovazzi 2017

Real boundary condition on Γ : $\mathbf{u} = \mathbf{u}_D$

Discrete bound. cond. on $\tilde{\Gamma}$: $\mathbf{u} + \nabla \mathbf{u} \cdot \mathbf{d} = \mathbf{u}_D$



Advantage

No cut integral

Difficulty

- Construction of d
- \mathbb{P}_k , Neumann conditions

Outline

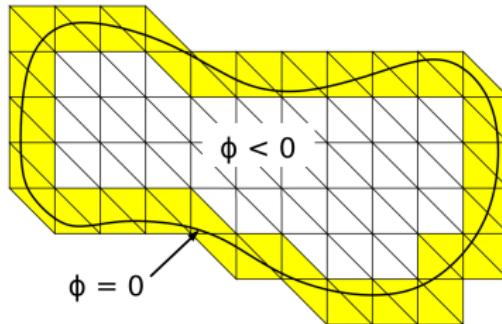
- The "standard" Finite Element Method
- Geometrical constraint on the mesh
- Previous Fictitious domain Methods
- ϕ -FEM
- Conclusions and perspectives

What is the idea of ϕ -FEM ?

Hypothesis :

Assume that Ω and Γ are given by a **level-set** function ϕ :

$$\Omega := \{\phi < 0\} \text{ and } \Gamma := \{\phi = 0\}.$$

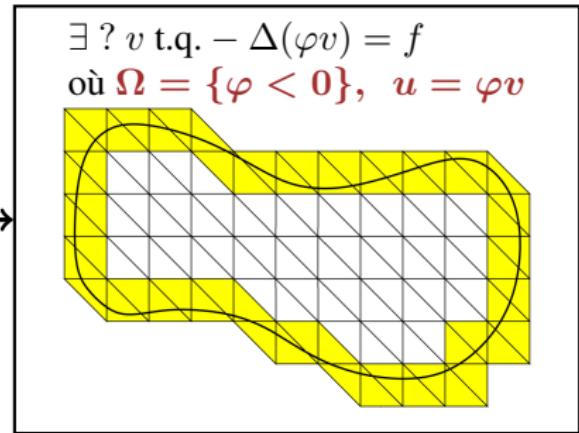
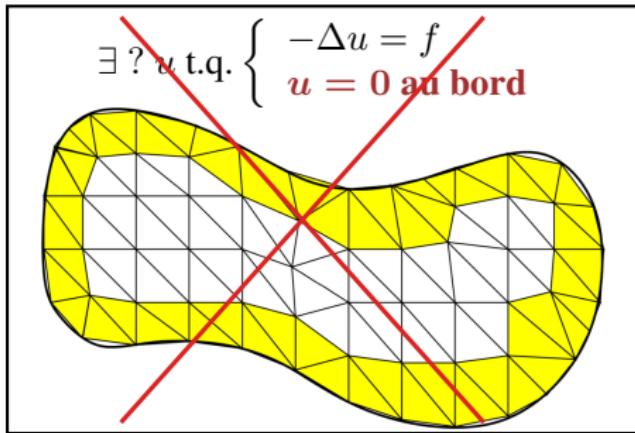


Idea of ϕ -FEM :

Include the Level-set function in the formulation
to take into account the boundary conditions

Formal approach

φ FEM for the Poisson Dirichlet Problem



ϕ -FEM : weak formulation

Hypothesis : Ω and Γ are given by a level-set function ϕ :

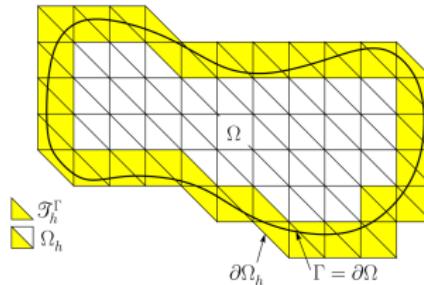
$$\Omega := \{\phi < 0\} \text{ and } \Gamma := \{\phi = 0\}.$$

ϕ -FEM formulation : Find v_h s.t.

$$\int_{\Omega_h} \nabla(\phi v_h) \cdot \nabla(\phi w_h) - \int_{\partial\Omega_h} \frac{\partial}{\partial n}(\phi v_h) \phi w_h + \text{Stab. Term} = \int_{\Omega_h} f \phi w_h \quad \forall w_h \in W_h,$$

where $W_h = \{\text{cont. piecewise pol. functions on a the mesh}\}$.

Then $u_h := v_h \phi$ as approximation of the initial problem.



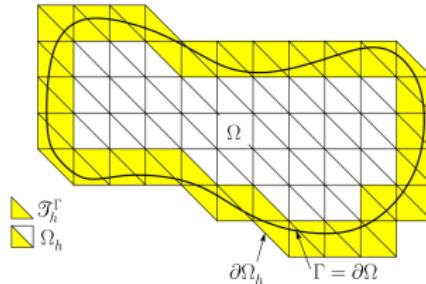
ϕ -FEM : weak formulation

Stabilization terms (to ensure a good conditioning number)

$$\begin{aligned} \sigma h \sum_{E \in \mathcal{F}_\Gamma} \int_E \left[\frac{\partial}{\partial n} (\phi_h v_h) \right] \left[\frac{\partial}{\partial n} (\phi_h w_h) \right] \\ + \sigma h^2 \sum_{T \in \mathcal{T}_h^\Gamma} \int_T (\Delta(\phi_h v_h) + f) \Delta(\phi_h w_h) \end{aligned}$$

where $[\cdot]$ is the jump on the interface E

$$\begin{cases} \mathcal{T}_h^\Gamma = \{T \in \mathcal{T}_h : T \cap \Gamma_h \neq \emptyset\} \quad (\Gamma_h = \{\phi_h = 0\}); \\ \mathcal{F}_\Gamma = \{E \text{ (an internal edge of } \mathcal{T}_h) \text{ such that } \exists T \in \mathcal{T}_h : \\ \quad T \cap \Gamma_h \neq \emptyset \text{ and } E \in \partial T\}. \end{cases}$$



ϕ -FEM : a priori error estimate

Theorem (D.-Lozinski 2020, optimal error)

Let u be the continuous solution and v_h the ϕ -FEM solution

Deformation error : $|u - u_h|_{1,\Omega} \leq Ch^k \|f\|_{k,\Omega_h}$

Displacement error : $\|u - u_h\|_{0,\Omega} \leq Ch^{k+1/2} \|f\|_{k,\Omega_h}$

with $C = C(\phi)$ and k the polynomial order.

Proposition (Optimal conditioning)

The finite element matrix of ϕ -FEM satisfies

$$\text{Cond}(\mathbf{A}) \leq Ch^{-2}.$$

ϕ -FEM : Simulation of the Poisson-Dirichlet problem

Domain Ω : circle of radius $\sqrt{2}/4$ centered at $(0.5, 0.5)$

Surrounding domain : $\mathcal{O} = (0, 1)^2$

Level-set function : $\phi(x, y) = (x - 1/2)^2 - (y - 1/2)^2 - 1/8$

Exact solution : $u(x, y) = \exp(x) \times \sin(2\pi y)$

Artificial external force : $f := -\Delta u$

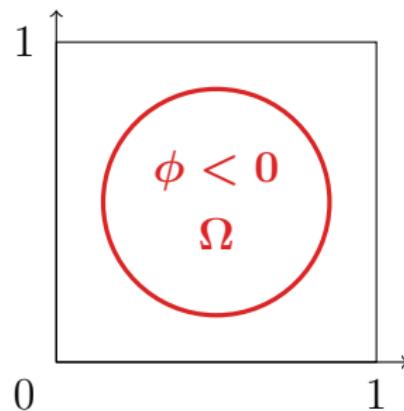
Boundary : $u_D := u(1 + \phi)$

Remark (non-homogeneous case)

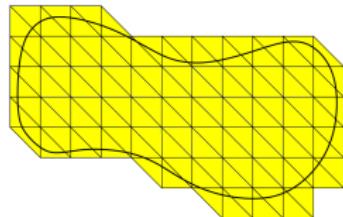
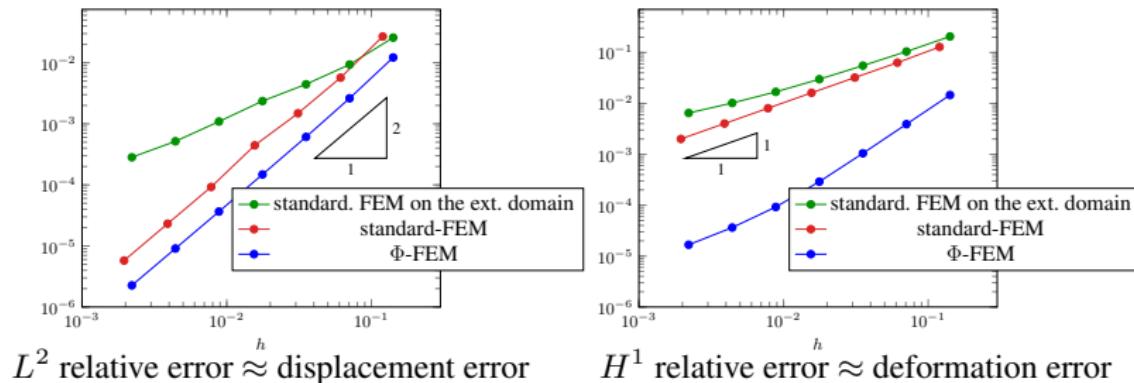
If $u := u_D$ on Γ ,

we replace $\phi_h v_h$ in the scheme

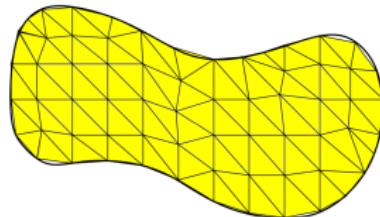
by $\phi_h v_h + u_D$ on Ω_h^Γ .



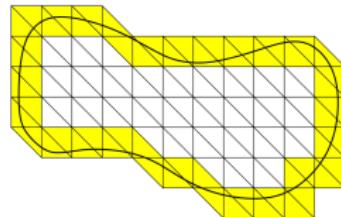
ϕ -FEM : Simulations in \mathbb{P}_1



standard FEM on the extended domain with a perturbation
(Dirichlet : $(1 + \phi)u$)



standard FEM

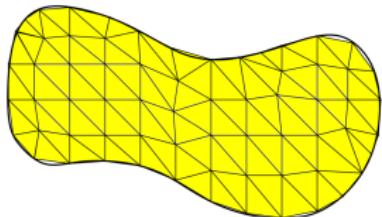


ϕ -FEM

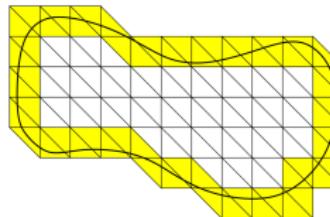
ϕ -FEM : Poisson-Dirichlet problem

Explanation of the numerical results

- Why better than standard FEM ?
 - Standard FEM : polygonal approximation of the boundary
 - ϕ -FEM : better approx. of the bound. with a levelset
- Numerical cost of ϕ -FEM :
 - Good point : small size of the FEM matrix
 - Bad point : quadrature more expensive
 - integral of polynomial product

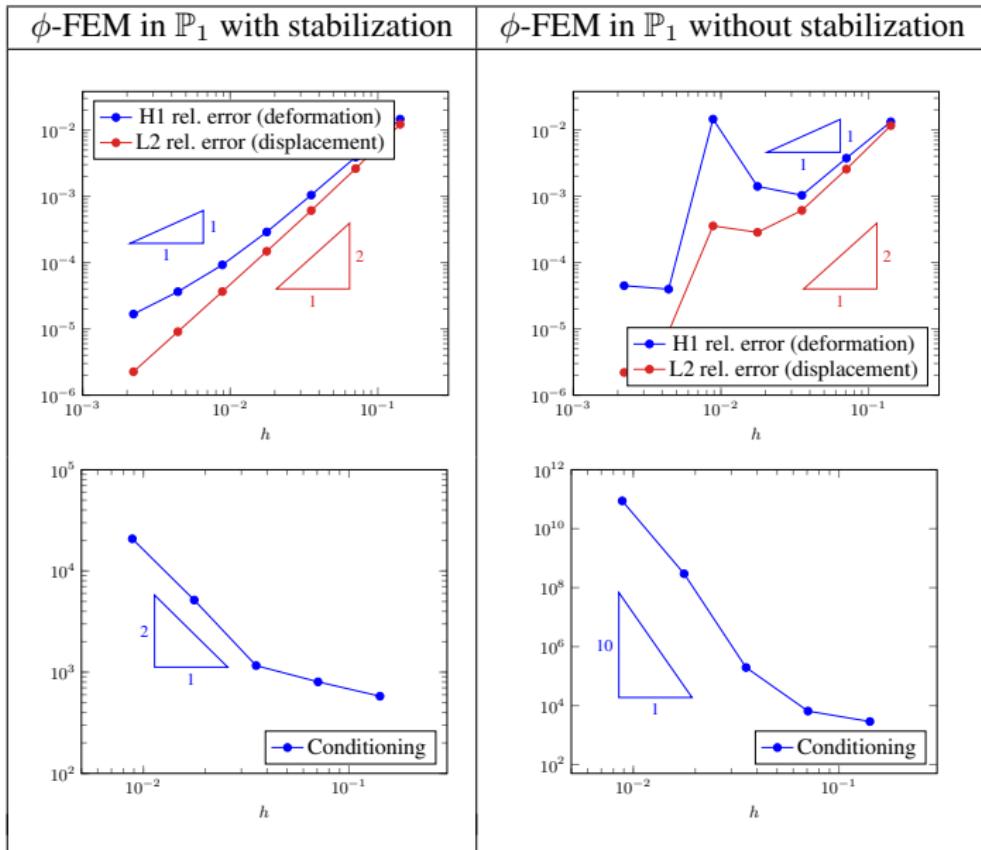


standard FEM



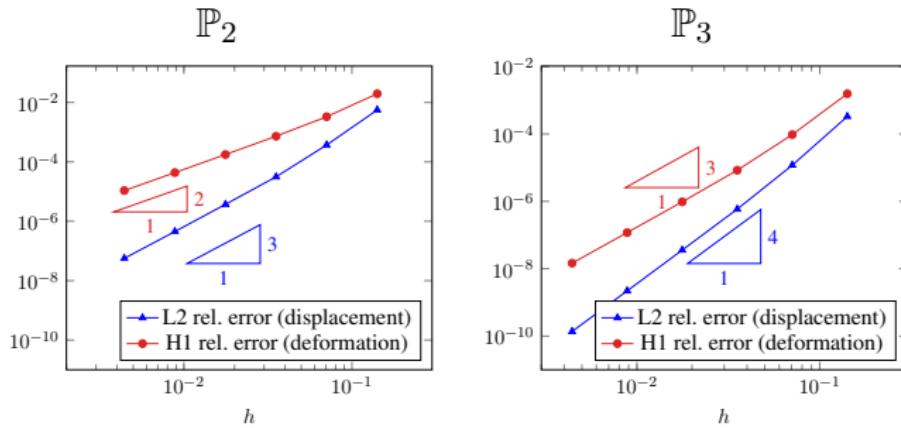
ϕ -FEM

ϕ -FEM : Stabilization and conditioning



ϕ -FEM : high polynomial order

Remarque : ϕ -FEM works high polynomial orders



Neumann boundary condition

Consider the **Poisson Neumann** problem :

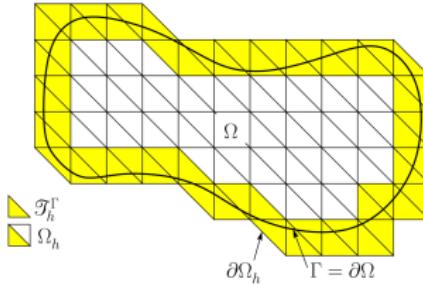
$$\begin{cases} -\Delta u + u = f, & \text{in } \Omega \\ \nabla u \cdot n = g, & \text{on } \partial\Omega \end{cases}$$

If

$$\boxed{\phi = 0 \quad \text{and} \quad \nabla\phi = n \text{ on } \Gamma},$$

the last system is (formally) equivalent to

$$\begin{cases} y = -\nabla u, & \text{in } \mathcal{T}_h^\Gamma \\ \mathbf{y} \cdot \nabla\phi = p\phi + g, & \text{in } \mathcal{T}_h^\Gamma \end{cases} \approx > \nabla u \cdot n = g \quad \text{on } \Gamma$$



Neumann boundary condition : weak formualtion

Find $(u_h, y_h, p_h) \in W_h^{(k)}$ such that for all $(v_h, z_h, q_h) \in W_h^{(k)}$

$$\begin{aligned}
& \int_{\Omega_h} \nabla u_h \cdot \nabla v_h + \int_{\Omega_h} u_h v_h + \int_{\partial\Omega_h} y_h \cdot n v_h \\
& \quad + \gamma_u \int_{\Omega_h^\Gamma} (y_h + \nabla u_h) \cdot (z_h + \nabla v_h) \\
& \quad + \frac{\gamma_p}{h^2} \int_{\Omega_h^\Gamma} (y \cdot \nabla \phi_h + \frac{1}{h} p_h \phi_h) (z \cdot \nabla \phi_h + \frac{1}{h} q_h \phi_h) \\
& \quad + \sigma h \int_{\Gamma_h^i} [\partial_n u] [\partial_n v_h] + \gamma_{div} \int_{\Omega_h^\Gamma} (\operatorname{div} y_h + u_h) (\operatorname{div} z_h + v_h) \\
= & \int_{\Omega_h} f v_h + \frac{\gamma_p}{h^2} \int_{\Omega_h^\Gamma} g (z \cdot \nabla \phi_h + \frac{1}{h} q_h \phi_h) + \gamma_{div} \int_{\Omega_h^\Gamma} f (\operatorname{div} z_h + v_h),
\end{aligned}$$

where φ_h is the Lagrange interpolation of ϕ of order l and

$$\begin{aligned}
W_h = \{ & (u_h, y_h, p_h) \in C^0(\Omega_h) \times C^0(\Omega_h^\Gamma)^d \times L^2(\Omega_h^\Gamma) : \\
& (u_h, y_h, p_h)_K \in \mathbb{P}_k \times (\mathbb{P}_k)^d \times \mathbb{P}_{k-1} \}
\end{aligned}$$

Neumann : optimal convergence

Theorem (D.-Lleras-Lozinski 2021, optimal error)

Let u be the continuous solution and u_h the ϕ -FEM solution

Deformation error : $|u - u_h|_{1,\Omega} \leq Ch^k(\|f\|_{k,\Omega_h} + \|g\|_{k+1,\Omega_h^\Gamma})$

Displacement error : $\|u - u_h\|_{0,\Omega} \leq Ch^{k+1/2}(\|f\|_{k,\Omega_h} + \|g\|_{k+1,\Omega_h^\Gamma})$

with $C = C(\phi)$ and k the polynomial order.

Proposition (Optimal conditioning)

The finite element matrix of ϕ -FEM satisfies

$$\text{Cond}(\mathbf{A}) \leq Ch^{-2}.$$

ϕ -FEM Neumann Poisson

- **Level set function :**
distance to the boundary
- **Exact solution :** $u(x, y) := \sin(x) \exp(y)$
- **Source term :** $f := -\Delta u + u$
- **Extrapolated Neumann boundary condition :**

$$\tilde{g} = \frac{\nabla u \cdot \nabla \phi}{|\nabla \phi|} + u \phi$$

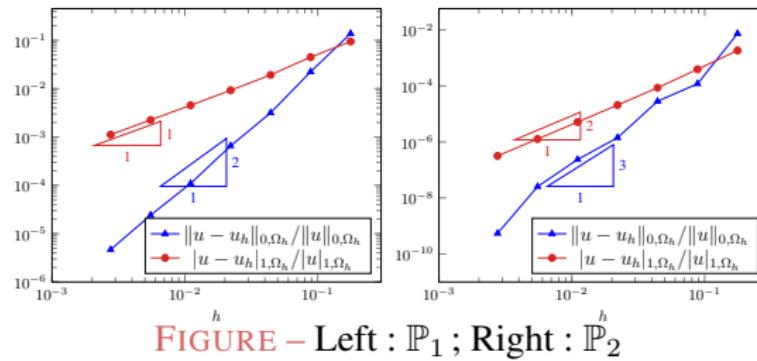
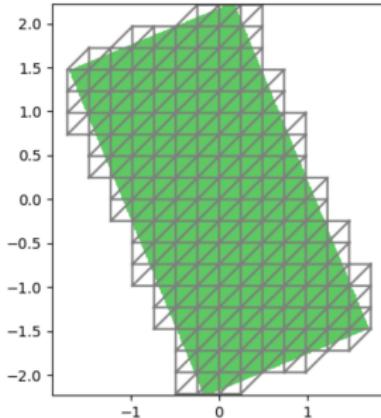


FIGURE – Left : \mathbb{P}_1 ; Right : \mathbb{P}_2

ϕ -FEM Neumann Poisson - Sensibility to cut cells

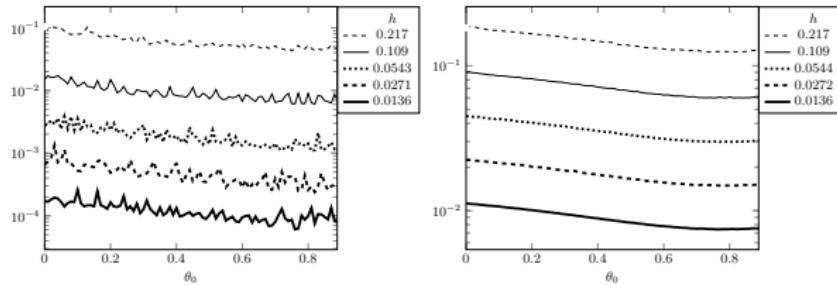


FIGURE – Sensitivity to the rotation in ϕ -FEM, $k = 1$ and $l = 3$. Left : L^2 relative error ; Right : H^1 relative error.

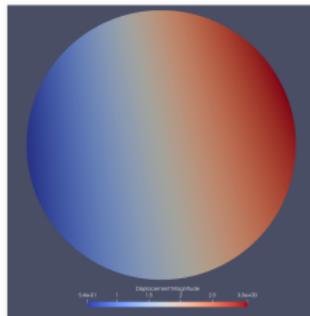
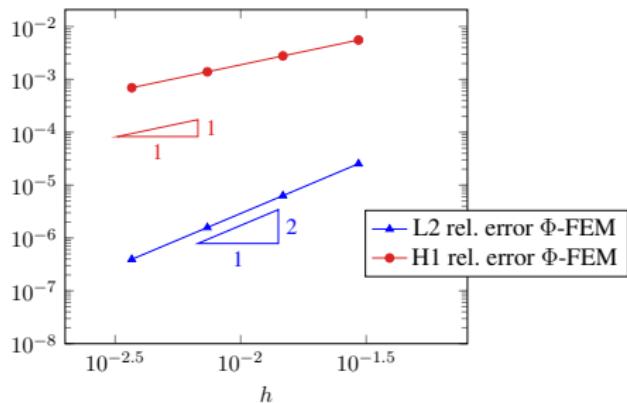
ϕ -FEM : Linear elasticity equation - Dirichlet

Model

$$\begin{cases} -\operatorname{div}(\sigma(u)) = f, & \text{in } \Omega, \\ u = u_D, & \text{on } \Gamma \end{cases}$$

ϕ -FEM model

$$-\operatorname{div}(\sigma(\phi v + u_D)) = f, \text{ in } \Omega,$$



[Duprez-Roussel]

Conclusion

Results

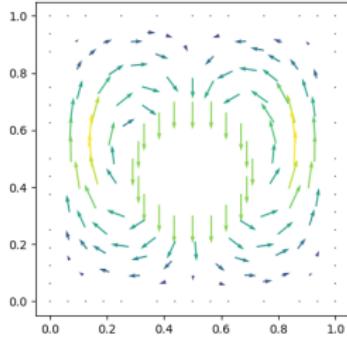
For the Poisson Dirichlet/Neumann problem :

- Optimal convergence
- Discrete problem well conditioned
- Simple implementation : standard shape functions
- Formulation available for any order of approximation

ϕ -FEM works for the elasticity equations

Work in progress

- Dynamic system
- Fluid structure
- Implementation of elastic models
in Sofa



Thanks for your attention !